

## PH 205: Mathematical methods of physics

### Problem Set 7

1. In this problem you will evaluate the Gaussian integral with complex parameters.

(a) First evaluate

$$\int_{-\infty}^{\infty} \exp [-\alpha(x-u)^2] dx,$$

$\alpha > 0$  and  $u$  is any complex number. Do this by converting the above integral to a complex integral along an appropriate contour and using the known result for

$$\int_{-\infty}^{\infty} \exp [-\alpha(x-x_0)^2] dx,$$

where  $x_0$  is any real number.

(b) Now, evaluate

$$\int_{-\infty}^{\infty} \exp [-\gamma x^2] dx,$$

where  $\gamma$  is a complex number. One again choose an appropriate contour and use the known result for the Gaussian integral. What are the constraints on  $\gamma$  for the above integral to be finite.

2. The Kramers-Kronig relations relate the real and imaginary parts of susceptibilities (also called response functions) in physics. Such a susceptibility  $\chi(\omega)$  is a function of angular frequency  $\omega$  and is of the form

$$\chi(\omega) = \chi_1(\omega) + \chi_2(\omega),$$

where  $\chi_1(\omega)$  and  $\chi_2(\omega)$  are the real and imaginary parts of the susceptibility.  $\chi(\omega)$  also has the property that it is analytic in the upper half plane and goes to zero faster than  $1/|\omega|$  as  $\omega \rightarrow \infty$ .

(a) Show that

$$\chi_1(\omega) = \frac{1}{\pi} \mathcal{P} \int_{-\infty}^{\infty} \frac{\chi_2(\omega')}{\omega' - \omega} d\omega',$$

$$\chi_2(\omega) = -\frac{1}{\pi} \mathcal{P} \int_{-\infty}^{\infty} \frac{\chi_1(\omega')}{\omega' - \omega} d\omega',$$

(b) The linear response of a time dependent quantity  $X(t)$  to a time dependent perturbation  $F(t)$  is given by

$$X(t) = \int_{-\infty}^{\infty} G(t-t') F(t') dt',$$

where  $G(t-t')$  is the Green's function which you have encountered before in problem 4 of problem set 5. In that problem you argued that  $G(\tau) = 0$  for  $\tau < 0$ . The susceptibility  $\chi(\omega)$  is the fourier transform of the Green's function and the above fact can be used to argue that  $\chi(\omega)$  is analytic in the upper half plane. Further, it can also be argued that for physical systems  $\chi(\omega)$  does indeed fall off faster than  $1/|\omega|$  as  $\omega \rightarrow \infty$ .  $X(t)$  and  $F(t)$  are real for for all  $t$  for physical systems and this implies that the Green's function is also a real-valued function. Use these facts to convert the integrals of part (a) to those from 0 to  $\infty$ .

3. In this problem you will calculate Green's functions (also called propagators) for the following operators, which are very important in physics, using the technique of Fourier transforms and contour integration:

- (a) The Helmholtz operator

$$\nabla^2 + m^2,$$

where  $m$  is a real number. Calculate the Green's function  $G(\mathbf{r} - \mathbf{r}')$  in 3 dimensions.

- (b) The D'Alembertian operator

$$\frac{1}{c} \frac{\partial^2}{\partial t^2} - \nabla^2.$$

Here  $t$  is time,  $c$ , the speed of light and  $\nabla^2$  is the usual Laplacian operator in three dimensions. Calculate the Green's function  $G(t - t'; \mathbf{r} - \mathbf{r}')$ .

- (c) The Schrödinger operator for a free particle in one dimension

$$i\hbar \frac{\partial}{\partial t} + \frac{\hbar^2}{2} \frac{\partial^2}{\partial x^2}.$$

Calculate the Green's function  $G(t - t'; x - x')$

- (d) The diffusion operator in one dimension

$$\frac{\partial}{\partial t} - D \frac{\partial^2}{\partial x^2},$$

where  $D > 0$  and is the diffusion constant. Calculate the Green's function  $G(t - t'; x - x')$

4. A very important technique in quantum statistical mechanics is that of Matsubara summation. In this problem you will perform such a sum using the techniques of contour integration. A typical Matsubara sum, which occurs in the physics of solids involving interactions between electrons and phonons at a temperature  $T$  is of the form

$$S = -\frac{1}{k_B T} \sum_n f(i\omega_n).$$

Here  $k_B$  is the Boltzmann constant and  $\omega_n = 2\pi n k_B T / \hbar$  where  $n$  is an integer, where  $\hbar$  is Planck's constant divided by  $2\pi$ . The  $\omega_n$ 's are called Matsubara frequencies and  $f(i\omega_n)$  is some (not necessarily analytic) function of  $\omega_n$ . The summation is over all integers  $n$ .

- (a) Show that the above sum can be evaluated by performing the integral

$$I = \oint_C \frac{dz}{2\pi i} f(z) n(z),$$

over a contour  $C$  which is a circle of radius  $R \rightarrow \infty$  centred at the origin. Here

$$n(z) = \left( \frac{1}{e^{\hbar z / k_B T} - 1} \right),$$

which you might recognize as the Bose distribution function (phonons are bosons) at zero chemical potential with  $\hbar z$  playing the role of energy. How is the sum related to the singularities of  $f(z)$  and their residues?

- (b) Now, consider a specific case of the function  $f$ , which is

$$f(i\omega_n) = \frac{2\Omega}{\omega_n^2 + \Omega^2} \frac{1}{ip_m + i\omega_n - E},$$

where  $\Omega$  and  $E$  are real numbers and  $p_m$  is also a Matsubara frequency. Evaluate

$$S = -\frac{1}{k_B T} \sum_n f(i\omega_n).$$

Note that the summation is over  $n$  and not  $m$ , which is some fixed integer. Physically, the parameters  $\Omega$  and  $E$  are the frequency of the phonon and the energy of the electron respectively though that is not important to performing the summation.

- (c) Suppose the summation had to be performed over all Matsubara frequencies of the form  $\omega_n = (2n + 1)\pi k_B T / \hbar$ . What would be an appropriate choice of  $n(z)$  in part (a)?

5. The Laplace transform of a function  $f(x)$  is defined as

$$F(s) = \int_0^\infty f(x) e^{-sx} dx.$$

- (a) Show that the Laplace transform of the function

$$f(x) = \frac{e^{ax}}{\sqrt{x^3}},$$

is

$$F(s) = -2\sqrt{\pi} \sqrt{s-a},$$

for  $s > a$ . You do not need to use contour integration for this: Use known results for Gaussian integrals and manipulate them appropriately.

- (b) The inverse Laplace transformation can be shown to be defined in terms of the Bromwich integral

$$f(x) = \frac{1}{2\pi i} \int_{\lambda-i\infty}^{\lambda+i\infty} F(s) e^{sx} ds,$$

where the integral is along a line parallel to the imaginary axis.  $\lambda$  is taken to be any real positive number larger than the real parts of all the singularities of  $F(s)e^{sx}$ . Evaluate the inverse Laplace transform of  $F(s)$  given in part (a) and verify that it is equal to the given function  $f(x)$ . To do this, you will need to choose an appropriate contour to evaluate the Bromwich integral. Note that the  $F(s)$  taken to be a function of a complex variable  $z$  has a branch point at  $z = a$ .