

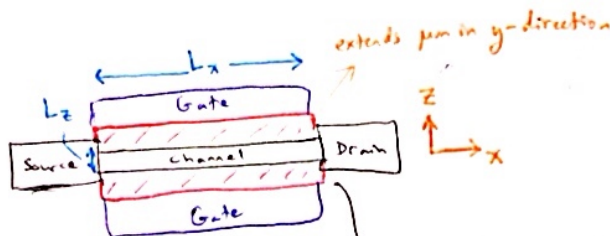
Lecture #5

Quantization Effects

Subbands

In a homogeneous solid (3D), electrons are free to move in all three directions and energy levels are classified: $E_b(k_x, k_y, k_z)$
bands, such as valence or conduction

What happens in this device?



Channel is a quantum well with e^- s free to move only in x - y plane

z -confinement modeled as a ring of circumference L_z , with periodic B.C. that restrict to allowed values of k_z given by: $k_z = \frac{p2\pi}{L_z}$

For a ring-shaped wire or nanotube this works

$$E_{b,p}(k_x, k_y) \approx E_b(k_x, k_y, k_z = \frac{p2\pi}{L_z})$$

$p \equiv$ subband index, $1, 2, 3, \dots$

The energy levels are now classified as:

$$E_{b,p}(k_x, k_y) \approx E_b(k_x, k_y, k_z = \frac{p2\pi}{L_z})$$

For the structure shown above, where the insulator layers act like ∞ potential wells

This is an approximation, but shows the following critical points

★ Quantum wells have 1D subbands, each having a 2D dispersion relation, $E(k_x, k_y)$.

★ How small does L_z need to be for something to be a 2D (quantum well)?

A: When it has experimentally observable effects.

→ Typically, the discrete energy levels of $k_z = \frac{p2\pi}{L_z}$ need to be $> k_B T$ (thermal energy)
Due to smoothing out of Fermi function

Example with parabolic $E-k$:

$$E(\vec{k}) \approx E_c + \frac{\hbar^2(k_x^2 + k_y^2 + k_z^2)}{2m_c}$$

Conduction band Edge

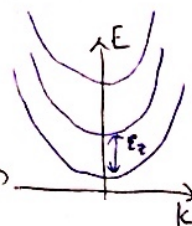
effective mass in conduction band

confinement in z gives rise to subbands:

$$E_p(k_x, k_y) \approx E_c + p^2 E_z + \frac{\hbar^2(k_x^2 + k_y^2)}{2m_c}$$

mass of an e^-

where: $E_z = \frac{\hbar^2 \pi^2}{2m_c L_z^2} = \frac{m_0}{m_c} \left(\frac{10 \text{ nm}}{L_z} \right)^2 \times 3.8 \text{ meV}$



Lecture # 5 (CONT)

Quantization Effects (CONT)

Subbands (CONT)

Example: Si, ETSOI : when is it 2D?

at R.T.

[IN class exercise]

$$m_c = 0.2 m_0, \quad k_B T \approx 26 \text{ meV @ R.T.}$$

$$\therefore \frac{1}{0.2} \left(\frac{10 \text{ nm}}{L_z} \right)^2 (3.8 \text{ meV}) > 26 \text{ meV}$$

$$\frac{1900 \text{ nm}}{L_z^2} > 26 \Rightarrow L_z < 8.5 \text{ nm}$$

OK, now what happens if we confine our channel in the y-direction as well to create a 1-D nanowire structure? Also called "quantum wires"



$$E_{n,p}(k_x) \approx E_c + n^2 E_y + p^2 E_z + \frac{\hbar^2 k_x^2}{2m_c}$$

$$E_y = \frac{\hbar^2 \pi^2}{2m_c L_y^2}$$

* Quantum wires (1D) have 2D subbands, each having a 1-D dispersion relation

Density of States

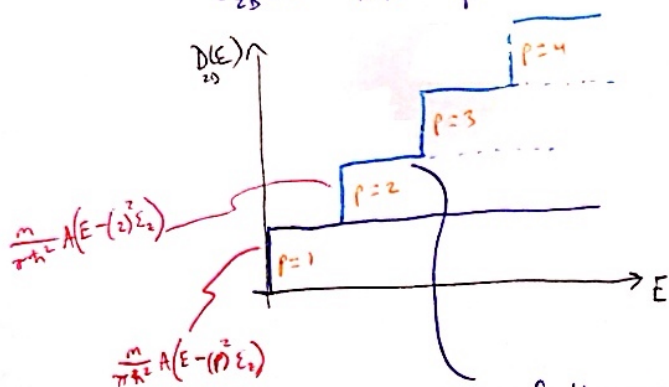
From Lecture 1 we found the 3D Dos : $D_{3D}(E) = \frac{m \sqrt{2mE}}{\pi^2 \hbar^3} \cdot V$ (includes spin)
Volume

Following the same approach for 2D gives: $D_{2D}(E) = \frac{m}{\pi \hbar^2} \cdot A$
Area

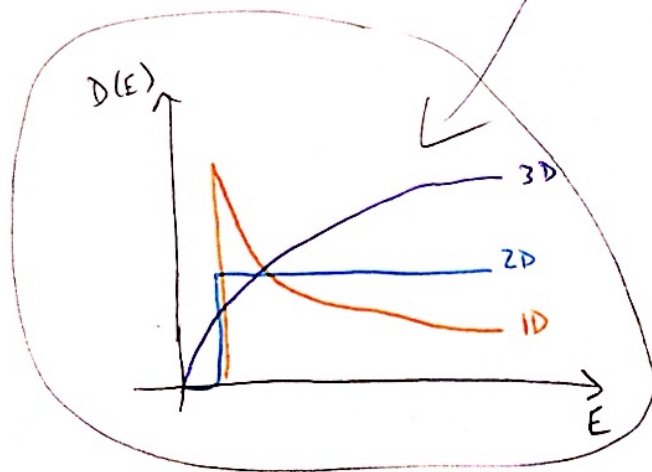
for 1D: $D_{1D}(E) = \frac{m}{\pi \hbar} (2mE)^{-1/2} \cdot L$
length

HOWEVER, the D_{2D} or D_{1D} are for only one subband! Must sum over all subbands:

$$D_{2D}(E) = \frac{m}{\pi \hbar^2} A \sum_p \psi(E - p^2 E_z)$$



Each of these subbands is called a "mode" for carrier transport



Lecture # 5 (cont)

Back to G_0

What is conductance of a material that is 1D, with perfect contacts?

Before, we said: $G_0 = \frac{2e^2}{h}$

Now, to consider the entire electronic structure:

$$G = M(E) G_0$$

of modes at E