Problem Bank 7: Partial Differential Equation

Kreyszig	Topics
Section	
12.1	Basic Concepts.
12.2-3	Wave Equation. Separating Variables.
	Use of Fourier Series.
12.6	Heat equation.

1 Basic Concepts.

1. Classification of PDEs

Classify the following equations in terms of its order, linearity and homogeneity (if the equation is linear).

(a) $u_t - u_{xx} + 1 = 0$

Solution: Second order, linear and non-homogeneous.

(b)
$$u_t - u_{xx} + xu = 0$$

Solution: Second order, linear and homogeneous.

(c)
$$u_t - u_{xxt} + uu_x = 0$$

Solution: Order 3 and non-linear.

(d)
$$u_{tt} - u_{xx} + x^2 = 0$$

Solution: Second order, linear and non-homogeneous.

(e)
$$\frac{u_x}{\sqrt{1+u_x^2}} + \frac{u_y}{\sqrt{1+u_y^2}} = 0$$

Solution: First order and non-linear.

(f)
$$u_t + u_{xxxx} + \sqrt{1+u} = 0$$

Solution: Order 4 and non-linear.

1 BASIC CONCEPTS.

2. Verify that for all pairs of differential functions f and g of one variable, u(x, y) = f(x)g(y) is a solution of the PDE $uu_{xy} = u_x u_y$.

Solution: First, compute u_x , u_y and u_{xy} :

$$u_x = g(y)f'(x)$$
$$u_y = f(x)g'(y)$$
$$u_{xy} = f'(x)g'(y)$$

Substituting into the PDE, we have

$$uu_{xy} = f(x)g(y)f'(x)g'(y) = u_x u_y$$

Hence, u(x, y) = f(x)g(y) is a solution of the PDE.

3. Boundary value problem

The Poisson's Equation is the non-homogeneous version of Laplace's Equation:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \rho(x, y) \tag{1}$$

Assume that $\rho(x, y) = 1$.

(a) Find the condition under which $u(x,y) = C_1 x^2 + C_2 y^2$ is a solution to the Poisson's Equation above.

Solution: First,
$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 2C_1 + 2C_2$$
. Since $\rho(x, y) = 1$, we have $C_1 + C_2 = \frac{1}{2}$.

(b) Suppose we also have the boundary condition

$$u(0,y) = \frac{y^2}{4}$$

Determine C_1 and C_2 .

Solution:
$$u(0,y) = C_1(0)^2 + C_2 y^2 = \frac{y^2}{4} \longrightarrow C_2 = \frac{1}{4}$$
. Hence, $C_1 = \frac{1}{2} - C_2 = \frac{1}{4}$.

(c) Find another solution that satisfies both the Poisson's Equation and the boundary condition.

Solution: Since $f(x,y) = \frac{1}{4}(x^2 + y^2)$ satisfies the Poisson's Equation, if we find any harmonic function g(x,y) such that $\frac{\partial^2 g}{\partial x^2} + \frac{\partial^2 g}{\partial y^2} = 0$, the sum of f and g will

continue to satisfy the Poisson's Equation. The boundary condition will be satisfied if g(0, y) = 0. One possible choice is therefore g(x, y) = x, which gives

$$u(x,y) = x + \frac{1}{4}(x^2 + y^2)$$

4. Initial–boundary value problem

Suppose a metal rod of length L has an initial temperature of $\sin\left(\frac{\pi}{L}x\right)$ and the temperatures at its left and right ends are both fixed at 0 degree Celsius. What would be the initial-boundary value problem that describes this scenario?

Solution:

PDE
$$u_t = c^2 u_{xx}, \quad 0 < x < L, \quad 0 < t < \infty$$

BCs $\begin{cases} u(0,t) = 0\\ u(L,t) = 0 \end{cases}, \quad 0 < t < \infty$
ICs $\begin{cases} u(x,0) = \sin\left(\frac{\pi}{L}x\right)\\ u_t(x,0) = 0 \end{cases}, \quad 0 \le x \le L$

5. Principle of Superposition

(a) Let \mathcal{D} be a unit disk centered at (0,0), i.e., \mathcal{D} includes all the points falling inside the unit circle $x^2 + y^2 = 1$. Suppose $f_1(x, y) = x^2 - y^2$ is a solution to the following boundary-value problem

$$\begin{cases} \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, & \text{for } (x, y) \in \mathcal{D} \\ u(x, y) = 2x^2 - 1, & \text{for } x^2 + y^2 = 1 \end{cases}$$

while $f_2(x, y) = x$ is a solution to the boundary value problem

$$\begin{cases} \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, & \text{for } (x, y) \in \mathcal{D} \\ u(x, y) = x, & \text{for } x^2 + y^2 = 1 \end{cases}$$

Solve the following boundary-value problem

$$\begin{cases} \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, & \text{for } (x, y) \in \mathcal{D} \\ u(x, y) = \frac{3}{4}(2x^2 - 1) + \frac{3}{5}x, & \text{for } x^2 + y^2 = 1 \end{cases}$$

Solution: By the principle of superposition,

$$u(x,y) = \frac{3}{4}(x^2 - y^2) + \frac{3}{5}x$$

will satisfy both the Laplace's Equation and the boundary condition.

1 BASIC CONCEPTS.

(b) Suppose $f_1(x,t) = \sin(\pi(x-ct))$ and $f_2(x,t) = \sin(3\pi(x+ct))$ both satisfy the wave equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

and the boundary conditions:

$$u(0,t) = 0, \quad u(1,t) = 0$$
 for all t

In addition, $u_1(x,t)$ and $u_2(x,t)$ satisfy the initial condition $u(x,0) = \sin(\pi x)$ and $u(x,0) = \sin(3\pi x)$, respectively. Find a solution that satisfies the initial condition $u(x,0) = \sin(\pi x) - \sin(3\pi x)$.

Solution: $u(x,t) = f_1(x,t) - f_2(x,t) = \sin(\pi(x-ct)) - \sin(3\pi(x+ct))$

6. Steady state solution

The heat equation for a rod with a constant internal heat source is described by the following PDE:

$$u_t = c^2 u_{xx} + 1, \quad 0 < x < 1$$

Suppose we fix the boundaries' temperatures by u(0,t) = 0 and u(1,t) = 1. Will the temperature u(x,t) converge to a constant temperature U(x) independent of time? If so, what is this constant temperature? HINT: Set $u_t = 0$.

Solution: Set $u_t = 0$, we have $u_{xx} = -\frac{1}{c^2}$. Integrating, we have

$$u(x) = -\frac{1}{c^2}\frac{x^2}{2} + Ax + B$$

Putting in the boundary conditions, we can find that B = 0 and $A = 1 + \frac{1}{2c^2}$. Hence the steady state temperature is given by

$$U(x) = -\frac{1}{2c^2}(x^2 - x) + x$$

- 7. Solve the following PDEs
 - (a) $u_{yy} = (\cosh x) yu;$
 - (b) $u_y = 2xyu;$

(c)
$$u_{xx} = 0, u_{yy} = 0,$$

where u = u(x, y).

Solution:

(a) We need to solve the partial differential equation

$$u_y = (\cosh x)yu.$$

Since no derivative of x occurs, the partial differential equation becomes

$$u' - (\cosh x)yu = 0 \Longrightarrow \frac{du}{u} = (\cosh x)ydy.$$

After integrating, we get

$$\ln u = \frac{1}{2}(\cosh x)y^2 + \ln A,$$

where we have

$$u = Ae^{\frac{1}{2}(\cosh x)y^2} \Longrightarrow u(x,y) = A(x)e^{\frac{1}{2}(\cosh x)y^2}$$

(b) We need to solve the partial differential equation

$$u_y = 2xyu.$$

Since no derivative of x occurs, the partial differential equation becomes

$$u' = 2xyu \Rightarrow \frac{u'}{u} = 2xy \Rightarrow \frac{du}{u} = 2xydy \Rightarrow \ln u = xy^2 + \ln A \Rightarrow u = Ae^{xy^2}.$$

Thus, we have $u(x,y) = A(x)e^{xy^2}$.

(c) Integrating the first PDE and the second PDE gives

 $u = c_1(y)x + c_2(y)$ and $u = c_3(x)y + c_4(x)$,

respectively. Equating these two functions gives

$$u = axy + bx + cy + k.$$

Alternatively, $u_{xx} = 0$ gives $u = c_1(y)x + c_2(y)$. Then from $u_{yy} = 0$, we get $u_{yy} = c_1''x + c_2'' = 0$, hence $c_1'' = 0, c_2'' = 0$, and by integration

$$c_1 = \alpha y + \beta, \quad c_2 = \gamma y + \delta$$

and by substitution in the previous expression

$$u = c_1 x + c_2 = \alpha x y + \beta x + \gamma y + \delta.$$

2 Wave Equation. Separating Variables. Use of Fourier Series.

1. D'Alembert Solution of the Wave Equation.

Using the D'Alembert solution, find the solution to the following initial-value problem

(a)

PDE
$$u_{tt} = c^2 u_{xx}, \quad -\infty < x < \infty, \quad 0 < t < \infty$$

ICs
$$\begin{cases} u(x,0) = e^{-x^2} \\ u_t(x,0) = 0 \end{cases}, \quad -\infty < x < \infty$$

(b)

PDE
$$u_{tt} = c^2 u_{xx}, \quad -\infty < x < \infty, \quad 0 < t < \infty$$

ICs
$$\begin{cases} u(x,0) = 0\\ u_t(x,0) = xe^{-x^2} \end{cases}, \quad -\infty < x < \infty$$

Solution: Applying the D'Alembert solutions directly, we have

(a)
$$u(x,t) = \frac{1}{2} (e^{-(x-ct)^2} + e^{-(x+ct)^2})$$

(b) $u(x,t) = \frac{1}{2c} \int_{x-ct}^{x+ct} u e^{-u^2} du = \frac{-1}{4c} \int_{x-ct}^{x+ct} e^{-u^2} d(-u^2) = \frac{1}{4c} \{e^{-(x-ct)^2} - e^{-(x+ct)^2}\}$

2. Wave equation and standing waves.

Find u(x,t) for the string of length L = 1 and $c^2 = 1$ when the string is fixed at the two ends (i.e., u(0,0) = u(L,0) = 0), the initial velocity is zero (i.e., $u_t(x,0) = 0$) and the initial shape (i.e., u(x,0) = 0) of the string is given by the specified function.

(a) $u(x,0) = kx(1-x^2)$

Solution:

It's given that $u_{tt} = u_{xx}$ and L = 1 with condition that

$$u_t(x,0) = 0$$
, and $u(x,0) = f(x) = kx(1-x^2)$.

The solution of $u_{tt} = u_{xx}$ is given by

$$u(x,t) = \sum_{n=1}^{\infty} B_n \cos n\pi t \sin n\pi x,$$

where $B_n = 2 \int_0^1 f(x) \sin n\pi x dx = 2k \int_0^1 x(1-x^2) \sin n\pi x dx.$

 $\int_0^1 x(1-x^2) \sin n\pi x dx$ can be obtained by using integration by parts

$$\int_0^1 x(1-x^2)\sin n\pi x dx = \frac{-6\pi n\cos n\pi}{\pi^4 n^4} = \frac{-6}{\pi^3 n^3} (-1)^n.$$

Thus we have
$$B_n = -\frac{12k}{\pi^3 n^3} (-1)^n$$
. Therefore we have the solution
$$u(x,t) = \sum_{n=1}^{\infty} B_n \cos n\pi t \sin n\pi x = \sum_{n=1}^{\infty} -\frac{12k}{\pi^3 n^3} (-1)^n \cos n\pi t \sin n\pi x$$

(b) u(x,0) = kx(1-x)

Solution:

It is given that $u_{tt} = u_{xx}$ and L = 1 with condition that

$$u_t(x,0) = 0$$
, and $u(x,0) = f(x) = kx(1-x)$. (2)

For L = 1, the solution of $u_{tt} = u_{xx}$ is given by

$$u(x,t) = \sum_{n=1}^{\infty} B_n \cos n\pi t \sin n\pi x,$$
(3)

where $B_n = 2 \int_0^1 kx(1-x) \sin n\pi x dx = 2k \int_0^1 x(1-x) \sin n\pi x dx$. After integration, we have

$$B_n = 2k \left[\left(-\frac{1}{n\pi} \right) (x(1-x)\cos n\pi x)_0^1 + \frac{1}{n\pi} \left\{ \frac{1}{n\pi} \left((1-2x)\sin n\pi x)_0^1 - \frac{2}{n^2\pi^2} (\cos n\pi x)_0^1 \right\} \right]$$
$$= 2k \left[\frac{2}{n^3\pi^3} \left(1 - (-1)^n \right) \right] = \frac{4k \left(1 - (-1)^n \right)}{n^3\pi^3}.$$

So we have

$$u(x,t) = \frac{4k}{\pi^3} \sum_{n=1}^{\infty} \frac{(1-(-1)^n)}{n^3} \cos n\pi t \sin n\pi x.$$
 (4)

(c) given by the following graph:



Solution: It is given that $u_{tt} = u_{xx}$ and L = 1 with conditions that $u_t(x,0) = 0$ and $u(x,0) = f(x) = \begin{cases} x & \text{if } 0 < x < 1/4 \\ \frac{1}{2} - x & \text{if } 1/4 < x < 1/2 \\ 0 & \text{if } 1/2 < x < 1 \end{cases}$

The solution of $u_{tt} = u_{xx}$ is given by $u(x,t) = \sum_{n=1}^{\infty} B_n \cos n\pi t \sin n\pi x,$ where $B_n = 2 \int_0^1 f(x) \sin n\pi x dx = 2 \left[\int_0^{1/4} x \sin n\pi x dx + \int_{1/4}^{1/2} \left(\frac{1}{2} - x\right) \sin n\pi x dx \right].$ After integration, we have $B_n = 2 \left[\left(-\frac{1}{n\pi} \right) (x \cos n\pi x)_0^{\frac{1}{4}} + \frac{1}{n^2 \pi^2} (\sin n\pi x)_0^{\frac{1}{4}} - \frac{1}{n\pi} \left\{ \left(\frac{1}{2} - x \right) \cos n\pi x \right\}_{\frac{1}{4}}^{\frac{1}{2}} - \frac{1}{n\pi} (\sin n\pi x)_{\frac{1}{4}}^{\frac{1}{2}} \right]$ $= \frac{2}{n^2 \pi^2} \left(2 \sin \left(\frac{n\pi}{4} \right) - \sin \left(\frac{n\pi}{2} \right) \right).$ Therefore, we have the solution as $u(x,t) = \frac{2}{\pi^2} \sum_{n=1}^{\infty} \frac{1}{n^2} \left(2 \sin \left(\frac{n\pi}{4} \right) - \sin \left(\frac{n\pi}{2} \right) \right) \cos n\pi t \sin n\pi x.$

Hint: If the initial velocity of the string is zero, the solution of the problem is given by

$$u(x,t) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) \cos\left(\frac{n\pi ct}{L}\right)$$

3. Method of separating variables

By the method of separating variables, write down the general solution for the equation

$$tu_t + u_{xx} + 2u = 0$$

under the boundary conditions

$$u(0,t) = u(\pi,t) = 0$$

You may proceed as follows:

(a) Search for separable solutions of the form $u(x,t) = f_1(t)f_2(x)$, to obtain two equations:

$$\begin{cases} \frac{tf'_1}{f_1} = -\lambda\\ \frac{f''_2 + 2f_2}{f_2} = -\lambda \end{cases}$$

- (b) When solving for $f_2(x)$, examine the following cases: $\lambda + 2 > 0$, $\lambda + 2 = 0$ and $\lambda + 2 < 0$. You should find that only in the first case can you find solutions that satisfy the boundary conditions.
- (c) Write down the general solution.

Solution:

• Let $u(x,t) = f_1(t)f_2(x)$, then we have

$$tf_2f_1' = f_1f_2'' + 2f_1f_2$$

• Separating variables and let the term on the two sides be $-\lambda$:

$$\frac{tf_1'}{f_1} = \frac{f_2'' + 2f_2}{f_2} = -\lambda$$

• The equation associated with t is

$$tf_1' = -\lambda f_1 \longrightarrow \frac{df_1}{f_1} = -\lambda \frac{dt}{t}$$
$$f_1(t) = Ae^{-\lambda \log t}$$

• The equation associated with x is

$$\frac{f_2''+2f_2}{f_2} = -\lambda \longrightarrow f_2'' + (\lambda+2)f_2 = 0$$

• For $\lambda + 2 > 0$, let $2 + \lambda = \beta^2$, we have $f_2'' + \beta^2 f_2 = 0$. So,

$$f_2(x) = C\cos\beta x + D\sin\beta x$$

Apply boundary conditions:

$$\begin{cases} f_2(0) = 0 \longrightarrow C = 0\\ f_2(\pi) = 0 \longrightarrow D \sin \beta \pi = 0 \end{cases}$$

which gives $\beta \pi = n\pi \longrightarrow \beta = n(n = 1, 2, 3, ...)$

- For $\lambda + 2 = 0$, the equation becomes $f_2'' = 0$, and $f_2(x) = Ax + B$. The boundary solutions are satisfied only when A = B = 0, thus the solution of this form is not useful.
- For $\lambda + 2 < 0$, the equation can be re-written as $f_2'' \gamma^2 f_2 = 0$, where γ is positive and $-\gamma^2 = \lambda + 2$. The solution in this case is $f_2 = Ae^{\gamma x} + Be^{-\gamma x}$. Again, the boundary solutions are satisfied only when A = B = 0, thus the solution of this form is not useful.
- The general solution is thus

$$u(x,t) = \sum_{n} (A_n \sin nx) e^{-\lambda \log t}$$

3 HEAT EQUATION.

3 Heat Equation.

1. Show that for a completely insulated bar described by the following boundary conditions:

$$u_x(0,t) = 0, \qquad u_x(L,t) = 0$$

and the following initial condition:

$$u(x,0) = f(x)$$

separation of variables gives the following solution

$$u(x,t) = A_0 + \sum_{n=1}^{\infty} A_n \cos \frac{n\pi x}{L} e^{-\left(\frac{cn\pi}{L}\right)^2 t}$$

where A_0 and A_n are given by

$$A_0 = \frac{1}{L} \int_0^L f(x) dx, \qquad A_n = \frac{2}{L} \int_0^L f(x) \cos \frac{n\pi x}{L} dx, \qquad n = 1, 2, \cdots$$

Solution: Since left hand side depends only on t and the right hand side only on x. So that both sides must be constant say k. For k = 0 and k > 0, we get only trivial solution i.e, u = 0, so we do not consider them.

Now for negative $k = ip^2, (p \neq 0)$, we have:

$$\frac{T'(t)}{c^2 T(t)} = \frac{X''(x)}{X(x)} = -p^2$$

$$X''(x) + p^2 X(x) = 0$$
 (1) and $T'(t) + c^2 p^2 T(t) = 0$ (2)

Now solution of the equation (1) is given by:

$$X(x) = A\cos px + B\sin bx$$

X'(0) = 0, so Bp = 0, and B = 0. Also, $X'(L) = 0 = -Ap \sin pl \Longrightarrow \sin pl = 0 \Longrightarrow p = \frac{n\pi}{L}$, $n = 1, 2, 3, \cdots$ Setting A = 1, we get the solution

$$X_n(x) = \cos(\frac{x\pi}{L})$$

Now, we solve equation (2) for $p = \frac{x\pi}{L}$:

$$T'(t) + c^2 \frac{n^2 \pi^2}{L^2} T(t) = 0$$

$$T'(t) + \lambda^2 n^2 T(t) = 0$$

where $\lambda_n = \frac{cn\pi}{L}$.

So its general solution is given by:

$$T_n = A_n e^{-\lambda n^2 t}, \qquad n = 0, 1, 2, \cdots$$

where A_n is a constant.

Hence, the functions

$$u_n(n.t) = X_n(x)T_n(t) = A_n \cos(\frac{nx\pi}{L})e^{-n^2t\lambda}$$

are the solutions of heat equation.

These are the eigenfunctions of the problem corresponding to the eigenvalues

$$\lambda_n = \frac{cn\pi}{L}$$

By the principle of superposition:

$$u(x,t) = \sum_{n=1}^{\infty} u_n(x,t) = \sum_{n=1}^{\infty} A_n \cos(\frac{nx\pi}{L}) e^{-n^2 t\lambda}$$

We have $\lambda_n = \frac{nc\pi}{L}$. Now,

$$u(x,0) = f(x) = \sum_{n=0}^{\infty} A_n \cos(\frac{nx\pi}{L})$$

where A_n is given by

$$A_n = \frac{2}{L} \int_0^L f(x) \cos(\frac{nx\pi}{L}) e^{\frac{-c^2 n^2 \pi^2 t}{L^2}}$$

Thus, the general solution is

$$u(x,t) = A_0 + \sum_{n=1}^{\infty} A_n \cos(\frac{nx\pi}{L}) e^{\frac{-c^2 n^2 \pi^2 t}{L^2}}$$

3 HEAT EQUATION.

2. Find the temperature in problem 1 above with $L = \pi$, c = 1, and $f(x) = 0.5 \cos 4x$.

Solution:

We need to solve the one-demensional equation:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

with the boundary conditions:

$$\begin{cases} u_x(0,t) = 0\\ u_x(\pi,t) = 0 \end{cases}$$

and the initial condition

$$u(x,0) = f(x) = 0.5 \cos 4x$$

We know that the general solution is given by

$$u(x,t) = \sum_{n=0}^{\infty} A_n \cos(nx) e^{-n^2} = A_0 + \sum_{n=1}^{\infty} A_n \cos(nx) e^{-n^2}$$

Now,

$$u(x,0) = f(x) = 0.5 \cos 4x = \sum_{n=1}^{\infty} A_n \cos(nx)$$

$$A_{n} = \frac{2}{\pi} \int_{0}^{\pi} f(x) \cos(nx) dx$$

= $\frac{2}{\pi} \int_{0}^{\pi} 0.5 \cos(4x) \cos(nx) dx$
= $\frac{1}{\pi} \int_{0}^{\pi} \cos(4x) \cos(nx) dx$
= $\begin{cases} \frac{1}{2} & \text{for} \quad n = 4\\ 0 & \text{for} \quad n \neq 4 \end{cases}$

That is, the only term that does not vanish is for n = 4. Thus, the solution is given by:

$$u(x,t) = \frac{1}{2}\cos(4x)e^{-16t}$$