#### LECTURE 10

## Algebraic formulas

I know it's easy to get lost in everything, so let me summarize our algebraic tools.

#### 1. Summary

**1.1.** The main players. There are two main algebraic players: First, we have the set of all deRham differential forms:

$$\Omega^{\bullet}_{deR}(M)$$

which is a ring. Its unit is the constant function  $1 \in \Omega^0_{deR}(M)$ . Its product is given by wedging together, pointwise.

The other player is

$$\Gamma(TM)$$

also known as vector fields. It is a Lie algebra.

**1.2.** Main properties. Skipping definitions, here are all the properties of the algebraic operations you need to know:

- d<sub>deR</sub> is a degree 1 derivation Ω<sup>•</sup><sub>deR</sub>(M) → Ω<sup>•</sup><sub>deR</sub>(M). It squares to zero.
  For any X ∈ Γ(TM), i<sub>X</sub> : Ω<sup>•</sup><sub>deR</sub>(M) → Ω<sup>•</sup><sub>deR</sub>(M) is a degree minus 1 derivation. It squares to zero.
- For any  $X, \mathcal{L}_X : \Omega^{\bullet}_{deR}(M) \to \Omega^{\bullet}_{deR}(M)$  is a degree zero derivation. It commutes with  $d_{deR}$ .

**1.3.** Useful formulas. Finally, here are the algebraic formulas you need to know.

• Cartan's magic formula,

$$\mathcal{L}_X := d_{deR} \circ \iota_X + \iota_X \circ d_{deR}$$

• How to compute the directional derivative of the values of a differential form:

$$\mathcal{L}_X(\alpha(Y_1,\ldots,Y_k)) = (\mathcal{L}_X\alpha)(Y_1,\ldots,Y_k) + \sum_{i=1}^k \omega(Y_1,\ldots,Y_{i-1},\mathcal{L}_XY_i,Y_{i+1},\ldots,Y_k).$$

• How the deRham derivative changes the values of a form:

$$d\alpha(Y_0, \dots, Y_k) = \sum_{i=0}^k (-1)^i Y_i(\alpha(Y_0, \dots, \hat{Y}_i, \dots, Y_k)) + \sum_{i < j} (-1)^{i+j} \omega([Y_i, Y_j], Y_0, \dots, \hat{Y}_i, \dots, \hat{Y}_j, \dots, Y_k)$$

That's it. Now I'll review the definitions of everything for your convenience.

## 2. The Lie derivative

Given any  $X \in \Gamma(TM)$ , one has maps

$$\mathcal{L}_X : \Gamma(TM) \to \Gamma(TM), \qquad \mathcal{L}_X : \Omega^k_{deR}(M) \to \Omega^k_{deR}(M).$$

which go by the same notation. To define this, we needed the theorem about ODEs: X defines a vector field, so you can flow along it. You evaluate  $\alpha$  at that point. Since flowing is a diffeomorphism, you can pull back the value of  $\alpha$  to the place you started before the flow.

$$(\Phi_t^X)^* \alpha(\Phi_t^X(x)) \in E_x.$$

Here, E is one of the vector bundles

$$TM, T^*M, \Lambda^k(T^*M),$$

and  $\alpha$  is a section of *E*. The Lie derivative says let's measure the rate of change of this vector with respect to *t*:

$$\mathcal{L}_X(\alpha)(x) := \lim_{t \to 0} \frac{(\Phi_t^X)^* \alpha(\Phi_t^X(x)) - \alpha(x)}{t}$$

#### 3. The Lie bracket

The Lie bracket was defined formally—if X and Y are two derivations from  $C^{\infty}(M)$  to itself, we saw that while  $X \circ Y$  may not be a derivation, the operation  $X \circ Y - Y \circ X$  is. So we defined

$$[X,Y] := X \circ Y - Y \circ X \in Der(C^{\infty}(M), C^{\infty}(M)).$$

Miraculously, this purely algebraic definition agrees with a more geometric operation:

$$\mathcal{L}_X Y = [X, Y].$$

We proved this last class. One can also see that

$$\mathcal{L}_X f = df(X) = X(f).$$

Making sure you can prove this is a very good way to check that you understand all the basic definitions.

#### 4. Differential forms eat tangent vectors

There is an isomorphism, for each k (and any finite-dimensional vector space V),

$$\Lambda^k(V^{\vee}) \to (\Lambda^k(V))^{\vee}$$

which sends an element of the form  $\alpha_1 \wedge \ldots \wedge \alpha_k$  to the linear map

 $v_1 \wedge \ldots \wedge v_k \mapsto \det(\alpha_i(v_j)).$ 

Since a differential form  $\alpha \in \Omega^k(M)$  defines an element

$$\alpha(x) \in \Lambda^k((T^*M)_x)$$

for any  $x \in M$ , the isomorphism tells us we can think of  $\alpha(x)$  as an element of

 $(\Lambda^k((TM)_x))^{\vee}$ 

i.e., as something that eats k tangent vectors to x and spits out a number. You can check that if  $f: M' \to M$  is a smooth map, then

 $((f^*\alpha)(x))(v_1,\ldots,v_k) = \alpha(f(x))(Df_x(v_1),\ldots,Df_x(v_k)).$ 

# 5. Contraction/interior multiplication

Given  $X \in \Gamma(TM)$ , we can define a map

$$\iota_X: \Omega^k(M) \to \Omega^{k-1}(M)$$

by sending

$$\alpha \mapsto \alpha(X, -, \dots, -).$$

This is defined using the isomorphism from the previous section. So for instance  $\iota_X(\alpha)$  eats in k-1 vector fields  $Y_1, \ldots, Y_{k-1}$  and spits out the function

$$\alpha(X, Y_1, \ldots, Y_{k-1}).$$

#### 6. deRham derivative $d_{deR}$

I write both  $d_{deR}$  and d depending on the context. They are the same thing. For every k, there is a map

$$d^k: \Omega^k(M) \to \Omega^{k+1}(M)$$

and I lazily call each of these maps d.

In local coordinates, this operation was defined as follows. Any k-form  $\alpha$  on  $U \subset \mathbb{R}^n$  can be written

$$\alpha = \sum_{I} f_{I} dx^{i_{1}} \wedge \ldots \wedge dx^{i_{k}}$$

where the index  $I = (i_1 < \ldots < i_k)$  runs through every sequence of k increasing integers between 1 and n. And  $f_I$  is a choice of smooth function on U. Then

$$d\alpha := \sum_{I} \sum_{j=1}^{n} \frac{\partial f_{I}}{\partial x^{j}} dx^{j} \wedge dx^{i_{1}} \wedge \ldots \wedge dx^{i_{k}}.$$

Though this is local, you can see that this defines a map of section  $\Omega^k(M) \to \Omega^{k+1}(M)$  by checking that this definition is compatible with the transition functions for the bundles  $\Lambda^k(T^*M)$  and  $\Lambda^{k+1}(T^*M)$ .

**Remark 10.1.** While formulas in local coordinates let you compute, those of you more formally-minded may ask for a more intrinsic characterization of the deRham derivative. There is a non-trivial theorem (see Conlon, Theorem 7.5.5 and its corollaries) asserting that

$$\Gamma(\Lambda^k(T^*M)) \cong \Lambda^k(\Gamma(T^*M)).$$

The lefthand side  $\Lambda$  takes exterior powers over  $\mathbb{R}$ ; the righthand side takes exterior powers over  $C^{\infty}(M)$ . So by the righthand side, demanding that d be a derivation means we only need specify it on 0-forms (functions) and 1-forms. This is what we did in class. (In fact, there is another universal property that characterizes d simply by declaring it on 0-forms; this will be an optional problem on a homework.)

#### 7. Some computational exercises

This is just a sanity check to make sure you know what you're doing. If you cannot do these, definitely come and talk to me and Phil. In what follows, everything is defined on  $\mathbb{R}^3$ .

- $\alpha = x^2 dx + y^2 dy + xy dz.$
- $\beta = xyzdx \wedge dy + zdy \wedge dz.$
- $\omega = xdy \wedge dx + xdx \wedge dy.$
- $\eta = dx \wedge dy \wedge dz$ .

• 
$$X = x \frac{\partial}{\partial x}$$
.

• 
$$Y = x^2 \frac{\partial}{\partial y}$$

(a) Show 
$$\omega = 0$$
.

- (b) Compute  $\iota_X \alpha$ .
- (c) Compute  $\iota_Y \beta$ .
- (d) Compute  $\iota_Y \eta$ .
- (e) Compute  $\alpha \wedge \beta$ .
- (f) Compute  $\beta \wedge \alpha$ .
- (g) Compute  $\mathcal{L}_X \alpha$ .

(h) Compute  $\mathcal{L}_Y\beta$ .

(i) Compute  $d\alpha$ .

(j) Compute  $d\beta$ .

(k) Compute [X, Y].

### 8. Solutions

(a) Show  $\omega = 0$ . Use that  $dy \wedge dx = -dx \wedge dy$ . More pedantically, at any point  $p \in \mathbb{R}^3$ , we have that

$$\omega(p) = x(p)dy|_p \wedge dx|_p + x(p)dx|_p \wedge dy|_p \in \Lambda^2((T^*\mathbb{R}^3)_p).$$

And this vector space was defined by quotienting out  $(T^*\mathbb{R}^3)_p \otimes (T^*\mathbb{R}^3)_p$  by imposing the relation  $v \otimes v = 0$ , which means that  $v \wedge w := [v \otimes w] = -[w \otimes v]$  $v] =: -w \wedge v$ . Long story short, we know that  $dy|_p \wedge dx|_p = -dx|_p \wedge dy|_p$ , so  $\omega(p) = 0$  for all p. This means that  $\omega$  is indeed the zero section of  $\Lambda^2 T^*\mathbb{R}^3$ .

Another, while equivalent, way to go about this: Note that

$$\Gamma(\Lambda^k(T^*\mathbb{R}^n)) \cong \Lambda^k(\Gamma(T^*\mathbb{R}^n))$$

where on the left, the exterior power is over  $\mathbb{R}$ , while on the right, it is over  $C^{\infty}(\mathbb{R}^n)$ . This is another way to reason out that  $dy \wedge dx = -dx \wedge dy$ . (b) **Compute**  $\iota_X \alpha$ .  $\iota_X$  is defined pointwise. That is,

$$(\iota_X \alpha)(p) := \iota_{X(p)} \alpha(p) = \alpha(p)(X(p)).$$

And  $dx(\frac{\partial}{\partial x}) = \frac{\partial x}{\partial x} = 1$ . Hence we have

$$\iota_X \alpha = x^2 dx(X) = x^2(x) = x^3.$$

One important point to note is that the value of  $\iota_X \alpha$  at a point p does not depend at all on how X or  $\alpha$  behave in a neighborhood of p, but only on their respective values at p.

One way to say this is that the map

$$\iota: \Gamma(TM) \times \Omega^k(M) \to \Omega^{k-1}(M)$$

is bilinear over  $C^{\infty}(M)$ , not just over  $\mathbb{R}$ . This is very useful and special. For instance, the Lie bracket is far from being bilinear over  $C^{\infty}(M)$ .

(c) **Compute**  $\iota_Y \beta$ . We see by linearity that

$$\iota_Y(\beta) = xyz\,\iota_Y(dx \wedge dy) + z\,\iota_Y(dy \wedge dz).$$

The easiest way to compute  $\iota$  is to pass through the isomorphism between  $\Lambda(V^{\vee})$  and  $(\Lambda(V))^{\vee}$ . Then

$$(dx \wedge dy)(Y, -) = dx(Y)dy(-) - dy(Y)dx(-)$$

and

$$(dy \wedge dz)(Y, -) = dy(Y)dz(-) - dz(Y)dy(-).$$
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Note that dx(Y) = dz(Y) = 0 while  $dy(Y) = x^2$ . Hence

$$\iota_Y(\beta) = xyz(-x^2)dx + z\,x^2dz = -x^3yzdx + x^2zdz.$$

Contracting by a vector field with only  $\frac{\partial}{\partial y}$  component, note, killed off the dy components of the differential form.

(d) **Compute**  $\iota_Y \eta$ . This gets more annoying if you use the isomorphism from  $\Lambda(V^{\vee})$  to  $(\Lambda(V))^{\vee}$  directly. Instead, recall the formula from class that

$$(\alpha \wedge \beta)(v_1, \dots, v_k) = \sum_{\pi \in k, l-\text{shuffles}} \alpha(v_{\pi(1)}, \dots, v_{\pi(k)}) \beta(v_{\pi(k_1)}, \dots, v_{\pi(k+l)}).$$

Since Y only has a  $\frac{\partial}{\partial y}$  component, let's take  $\alpha = dy$  and  $\beta = dz \wedge dx$ , so that  $\eta = \alpha \wedge \beta$ . Then, given any two vector fields  $V_2, V_3$ , we compute

$$(\alpha \wedge \beta)(Y, V_2, V_3)$$

by using the above formula. We see that the  $\beta$  factor is always zero if the shuffle  $\pi$  sends 1 to either 2 or 3. (After all,  $dz \wedge dx$  evaluated on two vectors is zero if one of the vectors has no  $\frac{\partial}{\partial x}$  or  $\frac{\partial}{\partial z}$  component.) So the only terms in the summation that are potentially non-zero are those that come from (1, 2)-shuffles that fix 1—i.e., with  $\pi(1) = 1$ . But the only such shuffle is the identity shuffle by definition of shuffle. Hence

$$(\alpha \wedge \beta)(Y, V_2, V_3) = \alpha(Y)\beta(V_2, V_3) = x^2\beta(V_2, V_3).$$

Since  $V_2, V_3$  were arbitrary, we conclude that

$$\iota_Y \eta = x^2 dz \wedge dx.$$

(e) **Compute**  $\alpha \wedge \beta$ . The answer is

$$(x^2z + x^2y^2z) \, dx \wedge dy \wedge dz.$$

(f) **Compute**  $\beta \wedge \alpha$ . The answer is

$$-(x^2z+x^2y^2z)\,dx\wedge dy\wedge dz.$$

(g) Compute  $\mathcal{L}_X \alpha$ . We compute

$$(8.1) \qquad \mathcal{L}_{X}(x^{2}dx + y^{2}dy + xydz) = \mathcal{L}_{X}(x^{2})dx + x^{2}\mathcal{L}_{X}(dx) + \mathcal{L}_{X}(y^{2})dy + y^{2}\mathcal{L}_{X}(dy) + \mathcal{L}_{X}(xz)dz + xy\mathcal{L}_{X}(dz)$$

$$(8.2) \qquad = \mathcal{L}_{X}(x^{2})dx + x^{2}d\mathcal{L}_{X}(x) + \mathcal{L}_{X}(y^{2})dy + y^{2}d\mathcal{L}_{X}(y) + \mathcal{L}_{X}(xz)dz + xyd\mathcal{L}_{X}(z)$$

$$(8.3) \qquad = 2x^{2}dx + x^{2}d(x) + 0dy + y^{2}d(0) + xzdz + xyd(0)$$

$$= 2x^{2}dx + x^{2}dx + xzdz$$

where in (8.1) we used that  $\mathcal{L}_X$  is a derivation for any X. In (8.2) we use that it commutes with d. Then in (8.3) we use that for functions,  $\mathcal{L}_X(f) = df(X) = X(f)$ , i.e., one computes the directional derivative of f in the direction of X. Since  $X = x \frac{\partial}{\partial x}$ , this entails computing  $\frac{\partial f}{\partial x}$ , then multiplying the result by x.

 $= 3x^2 dx + xz dz.$ 

(h) **Compute**  $\mathcal{L}_Y \beta$ . We compute

$$(8.4)$$

$$\mathcal{L}_{Y}(xyzdx \wedge dy + zdy \wedge dz) = \mathcal{L}_{Y}(xyz)dx \wedge dy + \mathcal{L}_{Y}(z)dy \wedge dz + xyz \mathcal{L}_{Y}(dx) \wedge dy + z \mathcal{L}_{Y}(dy) \wedge dz + xyz dx \wedge \mathcal{L}_{Y}(dy) + zdy \wedge \mathcal{L}_{Y}(dz)$$

$$(8.5) = \mathcal{L}_{Y}(xyz)dx \wedge dy + \mathcal{L}_{Y}(z)dy \wedge dz + xyz d(\mathcal{L}_{Y}(x)) \wedge dy + z d(\mathcal{L}_{Y}(y)) \wedge dz + xyz dx \wedge d(\mathcal{L}_{Y}(y)) + zdy \wedge d(\mathcal{L}_{Y}(z))$$

$$(8.6) = x^{3}zdx \wedge dy + 0 + 0 + 0 + z \cdot 2xdx \wedge dz + xyz dx \wedge 2xdx + 0$$

$$(8.7) = x^{3}zdx \wedge dy + 2xzdx \wedge dz.$$

Here, (8.4) again uses that  $\mathcal{L}_X$  is a derivation. (8.5) uses that  $\mathcal{L}_X$  commutes with d. (8.6) computes d of  $\mathcal{L}_Y(f)$  using that  $\mathcal{L}_Y(f) = Y(f)$ . (8.7) uses that  $dx \wedge dx = 0$ . (i) Compute  $d\alpha$ . Set  $x = x^1, y = x^2, z = x^3$ . Then

$$d\alpha = \sum_{i} \frac{\partial (x^{2})}{\partial x^{i}} dx^{i} \wedge dx + \sum_{i} \frac{\partial (y^{2})}{\partial x^{i}} dx^{i} \wedge dy + \sum_{i} \frac{\partial (xy)}{\partial x^{i}} dx^{i} \wedge dz$$
  
=  $2x \, dx \wedge dx + 2y \, dy \wedge dy + y \, dx \wedge dz + x \, dy \wedge dz$   
=  $0 + 0 + y \, dx \wedge dz + x \, dy \wedge dz$   
=  $y \, dx \wedge dz + x \, dy \wedge dz$ .

(j) Compute  $d\beta$ . Set  $x = x^1, y = x^2, z = x^3$ . Then

$$d\beta = \sum_{i} \frac{\partial (xyz)}{\partial x^{i}} dx^{i} \wedge dx \wedge dy + \sum_{i} \frac{\partial (z)}{\partial x^{i}} dx^{i} \wedge dy \wedge dz$$
  
=  $yz \, dx \wedge dx \wedge dy + xz \, dy \wedge dx \wedge dy + xy \, dz \wedge dx \wedge dy + dz \wedge dy \wedge dz$   
=  $0 + 0 + xy \, dz \wedge dx \wedge dy + 0$   
=  $xy \, dx \wedge dy \wedge dz$ .

(k) **Compute** [X, Y]. Recall that if

$$X = \sum X^i \frac{\partial}{\partial x^i}, \qquad Y = \sum Y^i \frac{\partial}{\partial x^i},$$

then

$$[X,Y] = \sum_{i,j} \left( X^i \frac{\partial Y^j}{\partial x^i} - Y^i \frac{\partial X^j}{\partial x^i} \right) \frac{\partial}{\partial x^j}.$$

So we have

$$\begin{split} [x\frac{\partial}{\partial x}, x^2\frac{\partial}{\partial y}] &= x\frac{\partial(x^2)}{\partial x}\frac{\partial}{\partial y} - x^2\frac{\partial(x)}{\partial y}\frac{\partial}{\partial x} \\ &= 2x^2\frac{\partial}{\partial y}. \end{split}$$