# Quantum Wells, Tunnels and Swings

### UP201

IISc Bangalore

## Of course, the questions..

- Item the second state is the second state of the second state o
- Whats with STM? (Scanning tunneling microscope)
- And quantum wells?
- Ramsauer effect alpha particles incident on inert gas atoms show "resonances" in transmission!
- Sound states? Metal-insulator transitions!

### The Great Barrier



- Model (What do you expect classically?)
- Scrödinger Recipe: Say energy state is E

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} = E\psi, \ |x| \ge a; \quad -\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} + V_0\psi = E\psi, \ |x| \le a$$

• Is  $E < V_0$  possible? What happens?

• Two length scales (*a* and 
$$\sqrt{\frac{\hbar^2}{2mV_0}}$$
)!!

## **Barrier Tunneling**

- Think of a state when electron is incident from left
- State to the left of the barrier x < -a: superposition of plane wave "traveling" to right and a "*reflected wave*"  $\psi_l(x) = e^{ikx} + Re^{-ikx}, \quad k = \sqrt{\frac{2mE}{\hbar^2}}$
- To the right of the barrier a wave traveling to right  $\psi_r(x) = Te^{ikx}$
- Restrict to case when  $E < V_0$
- The state in the barrier is  $\psi_b(x) = Ae^{qx} + Be^{-qx} (q = \sqrt{\frac{2m(V_0 E)}{\hbar^2}}) \text{does}$ *not* correspond to a wavelike state

## **Barrier Tunneling**

- How to determine *R*, *T*, *A*, *B*? Of interest: *T*!
- Use conditions

$$\psi_l(-a) = \psi_b(-a), \quad \psi'_l(-a) = \psi'_b(-a)$$
  
 $\psi_b(a) = \psi_r(a), \quad \psi'_b(a) = \psi'_r(a)$ 

four equations for four unknowns! All OK!

- Solution  $T = \frac{4ie^{2a(-ik+q)}kq}{(1-e^{4aq})q^2+2i(1+e^{4aq})kq-(1-e^{4aq})k^2}$
- The quantity  $|T|^2$  is probability of transmission...upshot  $|T|^2 \sim e^{-4aq}$ !
- For a given  $V_0$ , how does  $|T|^2$  depend on E, for various values of  $a/\sqrt{\frac{\hbar^2}{2mV_0}}$ ?

# **Barrier Tunneling**



- For small *a* the electron does not feel the barrier, for large *a* there is very little transmission
- Note "resonances" for large values of *a*, when *E* > *V*<sub>0</sub> (not considered, so far)

#### Tunneling in Real Life

- Alpha particles: Held to nucleus by *very strong* nuclear forces. Barrier height can be of many eVs!! Very high kinetic energy (confinement in  $10^{-15}$ m)!!  $|T|^2$  is quite small (in fact, can be shown to be  $\sim e^{-90}$ (!!!!)), the number of "attempts" to jump out are large (of the order of  $10^{21}$  per second)...probaility of decay is  $10^{-11}$  per year!!
- Ohmic contacts between metals and semiconductors: Need to have linear *I V* characteristics for contacts...problem...contact barrier! Solution: Try to make *a* as small as possible by heavily doping semiconductor near contact...depletion zone size comes down and electrons tunnel across!

### Tunneling in Real Life

- Field emission devices...pull out electrons using voltage–flat panel displays
- STM...maintain a constant tunneling current (at constant voltage) by adjusting distance...maps out topography of surface!



Well, What Now?...

$$V(x) = 0$$

$$V(x) = -V_{0}$$

$$V(x) = 0$$

- Model (What do you expect classically?)
- Scrödinger Recipe: Say energy state is E

$$-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} \quad = \quad E\psi, \quad |x| \ge a; \quad -\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} - V_0\psi = E\psi, \quad |x| \le a$$

• What happens when E > 0? Is E < 0 possible?

### Quantum Well

- Think of a state when electron is incident from left
- State to the left of the well x < -a: superposition of plane wave "traveling" to right and a "*reflected wave*"  $\psi_l(x) = e^{ikx} + Re^{-ikx}, \quad k = \sqrt{\frac{2mE}{\hbar^2}}$
- To the right of the well a wave traveling to right  $\psi_r(x) = Te^{ikx}$
- Restrict to case when *E* > 0
- The state in the well is  $\psi_w(x) = Ae^{iqx} + Be^{-iqx}$   $(q = \sqrt{\frac{2m(E+V_0)}{\hbar^2}}) does$  correspond to a wavelike state of shorter wavelength!
- Solve for *T* in exactly the same way as the barrier problem!

# Quantum Well



- There are particular values of *k* at which T = 1... "Resonances"...Why?
- Roughly explains the Ramsauer effect!
- What happens if E < 0?

### Quantum Well – Bound States



• States with  $-V_0 < E < 0$  are possible! "Bound states"!

• Wavefunction 
$$(q = \sqrt{\frac{2m|E|}{\hbar^2}}, k = \sqrt{\frac{2m(E-V_0)}{\hbar^2}})$$
  
 $\psi_l(x) = Ae^{qx}, \quad \psi_w(x) = Be^{ikx} + Ce^{-ikx}, \quad \psi_r(x) = De^{-qx}$ 

- Main idea: Wavefunction decay exponentially outside the box!
- How to determine allowed *Es*?

### Quantum Well – Bound States

• Use exactly same ideas as before

$$\psi_l(-a)=\psi_w(-a), \qquad \psi_l'(-a)=\psi_w'(-a) \ \psi_w(a)=\psi_r(a), \qquad \psi_w'(a)=\psi_r'(a)$$

- Main point: *Homogeneous equations for* A, B, C, D
- Condition of non-triviality gives possible energy states

$$\sqrt{\frac{|E|}{V_0 - |E|}} = \tan\left(\sqrt{\frac{2m(V_0 - |E|)a^2}{\hbar^2}}\right) \quad n \text{ even}$$
$$\sqrt{\frac{|E|}{V_0 - |E|}} = -\cot\left(\sqrt{\frac{2m(V_0 - |E|)a^2}{\hbar^2}}\right) \quad n \text{ odd}$$

### Bound States...Big Deal?

- Whats the big deal?
- The "simplest" model of an atom!!
- Bound state are always possible in 1D (and 2D..I am not sure about this) for "attractive" potentials
- Quantum wells are used for many purposes... Infrared sensor, lasers etc. etc..

### Bound States...Big Deal?

• In 3D there are attractive potentials for which there are no bound states! For example,  $V(r) \sim \frac{e^{-\xi r}}{r}$ , there is no bound state for  $\xi > \xi_c$ ...crucial to understand metal-insulator transitions (Mott)





- Tunneling : Transmission probability goes as  $e^{-2\ell q}$  where  $\ell$  is the size scale of the barrier
- Quantum wells: "Resonances"
- Bound states: Again useful for many things!

# Why Harmonic Oscillator?

• What is with specific heat of solids?



- Molecular vibrations?
- Light itself...really?

## 1D Harmonic Oscillator

• Hamiltonian ( $\omega$  – "natural frequency")

$$H = \frac{P^2}{2m} + \frac{1}{2}m\omega^2 X^2$$

• What are the energy eigenvalues and eigenstates? Need to solve

$$H\psi = E\psi$$

• Are there bound states? Extended states? Both?

## 1D Harmonic Oscillator

- What can we say without solving anything?
- Well, what do you expect classically? Any energy *E* is possible. Particle will be confined to  $|x| \le \sqrt{\frac{2E}{m\omega^2}}$
- Looks like we will have only bound states!
- Clearly,  $\langle X \rangle = 0!$

### 1D Harmonic Oscillator

• Since, 
$$\langle X \rangle = 0$$
,  $\langle V(X) \rangle = \frac{1}{2}m\omega^2 \Delta x^2$ 

- Since bound state,  $\langle P \rangle = 0$  and  $\langle \frac{P^2}{2m} \rangle = \frac{\Delta p^2}{2m}$
- From uncertainty relation  $\Delta p = \frac{\hbar}{\Delta x}$
- Total energy  $E = \frac{\hbar^2}{2m\Delta x^2} + \frac{1}{2}m\omega^2\Delta x^2$
- We can show this energy is minimum if  $\Delta x = \sqrt{\frac{\hbar}{m\omega}}$ , for this value of  $\Delta x$ ,  $E \approx \hbar \omega$ !
- Thus, uncertainty principle tells us that the minimum energy level of the harmonic oscillator must have energy of the order of ħω!
- Contrast this with classical result!

# Solution of Quantum Swing

- Solution of  $-\frac{\hbar^2}{2m}\frac{\partial^2\psi}{\partial x^2} + \frac{1}{2}m\omega^2 x^2\psi = E\psi$  is a bit technical
- Allowed energy levels  $E_n = (n + \frac{1}{2})\hbar\omega$ , n = 0, 1, 2...
- Lowest energy level  $\frac{1}{2}\hbar\omega!$  Very close to our estimate!
- Levels are equally spaced
- Energy eigenstates  $\psi_n(x) = C_n H_n(\alpha x) e^{-\frac{1}{2}\alpha^2 x^2}$  with  $\alpha = \sqrt{\frac{m\omega}{\hbar}}$ ,  $C_n = \sqrt{\frac{\alpha}{\sqrt{\pi 2^n n!}}}$
- $H_n(\xi)$  are the Hermite polynomials which satisfy the differential equation  $H''_n(\xi) 2\xi H'_n(\xi) + 2nH_n = 0$

### Hermite Polynomials...What?

•  $H_n(\xi)$  are the Hermite polynomials which satisfy the differential equation  $H''_{n}(\xi) - 2\xi H'_{n}(\xi) + 2nH_{n} = 0$  $H_0(\xi) = 1$  $H_1(\xi) = 2\xi$  $H_{2}(\xi) = 4\xi^{2} - 2$  $H_{2}(\xi) = 8\xi^{3} - 12\xi$  $H_4(\xi) = 16\xi^4 - 48\xi^2 + 12$  $H_{5}(\xi) = 32\xi^{5} - 160\xi^{3} + 120\xi$  $H_{6}(\xi) = 64\xi^{6} - 480\xi^{4} + 720\xi^{2} - 120$  $H_7(\xi) = 128\xi^7 - 1344\xi^5 + 3360\xi^3 - 1680\xi$  $H_{8}(\xi) = 256\xi^{8} - 3584\xi^{6} + 13440\xi^{4} - 13440\xi^{2} + 1680$  $H_{9}(\xi) = 512\xi^{9} - 9216\xi^{7} + 48384\xi^{5} - 80640\xi^{3} + 30240\xi$ 

(Pauling and Wilson)

#### What does the ground state look like?



(Pauling and Wilson)

#### What do other states look like?



- Much cleaner way to derive energy levels...play God with Driac notation!
- Intrinsic length scale  $x_o = \sqrt{\frac{\hbar}{m\omega}}$ , momentum scale  $p_o = \sqrt{\hbar m\omega}$
- Define *nondimensional operators*  $\hat{X} = \frac{X}{x_o}$  and  $\hat{P} = \frac{P}{p_o}$
- Hamiltonian  $H = \hbar \omega \left(\frac{1}{2}\hat{P}^2 + \frac{1}{2}\hat{X}^2\right)$
- Commutation relation  $[\hat{X}, \hat{P}] = i$

- *Define* annihilation operator  $a = \frac{\hat{X} + i\hat{P}}{\sqrt{2}}$
- Creation operator: Hermitian conjugate  $a^{\dagger} = \frac{\hat{X} i\hat{P}}{\sqrt{2}}$

• Clearly 
$$\hat{X} = \frac{a^{\dagger} + a}{\sqrt{2}}$$
 and  $\hat{P} = \frac{a - a^{\dagger}}{\sqrt{2}i}$ 

- Commutation relation  $[\hat{X}, \hat{P}] = i$  implies  $[a, a^{\dagger}] = 1$
- Hamiltonian  $H = \hbar \omega \left(\frac{1}{2}\hat{P}^2 + \frac{1}{2}\hat{X}^2\right) = \hbar \omega \left(a^{\dagger}a + \frac{1}{2}\right)!!$
- Clearly, *H* and  $a^{\dagger}a$  have the same eigenvectors! Thus, if  $a^{\dagger}a|\alpha\rangle = \alpha |\alpha\rangle$ ,  $H|\alpha\rangle = \hbar\omega(\alpha + \frac{1}{2})|\alpha\rangle$
- Anticipating, lets us call  $a^{\dagger}a = N$ , a Hermitian operator

• Suppose, 
$$N|\alpha\rangle = \alpha |\alpha\rangle$$
, what is  $a|\alpha\rangle$ ?

- $Na|\alpha\rangle = a^{\dagger}aa|\alpha\rangle = (aa^{\dagger} 1)a|\alpha\rangle = (\alpha 1)a|\alpha\rangle!$  This implies  $a|\alpha\rangle = C|\alpha 1\rangle$ , with  $C = \sqrt{\alpha}!$  Annihilation operator!
- Similarly,  $a^{\dagger} | \alpha \rangle = \sqrt{\alpha + 1} | \alpha + 1 \rangle$ ! Creation operator!
- "Clearly": a<sup>m</sup>|α⟩ = √α(α − 1)(α − 2)...(α − (m − 1)|α − m⟩...this necessarily implies that α has to be a nonnegative integer! Thus α = 0, 1, 2...etc. We have shown that the eigenvalues of N are nonnegative integers i. e., N|n⟩ = n|n⟩!

• 
$$H = \hbar\omega \left(N + \frac{1}{2}\right), H|n\rangle = \underbrace{\hbar\omega \left(n + \frac{1}{2}\right)}_{E_n}|n\rangle$$

- The ground state is  $|0\rangle$ ...called "vacuum" state
- Any higher state can be generated from the ground state  $|n\rangle = \frac{(a^{T})^n}{\sqrt{n!}}|0\rangle$
- Modern theory of solids is formulated using creation and annihilation operators, the number n being interpreted as the "number of particles"
- The above are called Bose operators (applicable to phonons); for electrons there are Fermi operators which satisfy *anti-commutation relations*!

#### Phonons in solids

• Quantized lattice vibrations! Phonon frequency depends on wavelength



Ashcroft and Mermin

#### Phonons, and other things

- Quantum mechanics explains specific heat anomaly (if k<sub>b</sub>T ≪ ħω, C<sub>v</sub> ~ e<sup>-ħω/k<sub>b</sub>T</sup>, if k<sub>b</sub>T ≫ ħω, C<sub>v</sub> ~ k<sub>b</sub> (Dulong-Petit)) Quantum alacrity!
- Can be used to understand molecular vibrations
- Quantum theory of light is developed in "exact analogy" with the theory of phonons...can show that light of frequency has energy levels  $\hbar\omega(n+\frac{1}{2})$  where *n* is the number of photons...if you have large *n* you will have very intense light
- More interestingly this will tell you that "vacuum" in quantum mechanics is a VERY BUSY place!

### Summary – Quantum Swing

- Energy levels go as  $\left(n+\frac{1}{2}\right)\hbar\omega$
- There is a very nice way to do this using creation and annihilation operators
- Explains many things about materials example specific heat anomaly
- Quantum swing is a very basic problem in physics, it will arise in one form or other in *many* problems