

PH 6297: Classical Field Theory

Homework Set 1

August 21, 2017

1. In class we reviewed Lorentz boost transformations in the case when one inertial frame (S') was moving wrt to the another, S along the common x-axis:

$$\begin{aligned}x'^1 &= \gamma(x^1 - \beta x^0), \\x'^{(2,3)} &= x^{(2,3)}, \\x'^0 &= \gamma(x^0 - \beta x^1).\end{aligned}$$

A. Show that one can generalize these to derive the form of the Lorentz transformations for the case when S' is moving away from S with a velocity not along x-axis, but in an arbitrary direction, i.e. $\mathbf{V} = V^1\hat{\mathbf{e}}_1 + V^2\hat{\mathbf{e}}_2 + V^3\hat{\mathbf{e}}_3$ by completing the hints provided in class (or independent of that):

$$\begin{aligned}\mathbf{x}' &= \mathbf{x} + \frac{\gamma - 1}{\beta^2}(\mathbf{x} \cdot \boldsymbol{\beta})\boldsymbol{\beta} - \gamma\boldsymbol{\beta}x^0, \\x'^0 &= \gamma(x^0 - \boldsymbol{\beta} \cdot \mathbf{x}),\end{aligned}$$

where, $\boldsymbol{\beta} \equiv \mathbf{V}/c$ and $\beta^2 \equiv \boldsymbol{\beta} \cdot \boldsymbol{\beta}$.

B. Write the above pair of equations first in component form, i.e. $x'^i = x^i + \dots$, $x'^0 = \gamma x^0 + \dots$ and then in matrix form,

$$x' = \Lambda x$$

indicating clearly the entries, $\Lambda^\mu{}_\nu$ in terms of γ, β^i .

Solution

A. In class we saw how to generalize the expression for Lorentz transformation from one inertial frame, S to another frame, S' moving with a velocity $\mathbf{v} = v\hat{\mathbf{e}}_1$ along the common x -axis to the case when S' is moving with a velocity, $\mathbf{v} = v_i \hat{\mathbf{e}}_i$, in an arbitrary direction by

using the following recipe,

$$\begin{aligned}x'^0 &= \gamma (x^0 - \beta x^{\parallel}), \\x'^{\parallel} &= \gamma (x^{\parallel} - \beta x^0), \\x'^{\perp} &= \mathbf{x}^{\perp},\end{aligned}$$

where $x^{\parallel} \equiv \mathbf{x} \cdot \hat{\boldsymbol{\beta}}$, $x'^{\parallel} = \mathbf{x}' \cdot \hat{\boldsymbol{\beta}}$ denote the component of position vector along the/ parallel to the velocity direction, while $\mathbf{x}^{\perp} \equiv \mathbf{x} - (\mathbf{x} \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}}$ and $\mathbf{x}'^{\perp} = \mathbf{x}' - (\mathbf{x}' \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}}$ denote the part of the position vector perpendicular to the velocity, $\boldsymbol{\beta} = \frac{\mathbf{v}}{c}$. Thus we have, the second equation,

$$\begin{aligned}x'^{\parallel} &= \gamma (x^{\parallel} - \beta x^0) \\ \implies \mathbf{x}' \cdot \hat{\boldsymbol{\beta}} &= \gamma (\mathbf{x} \cdot \hat{\boldsymbol{\beta}} - \beta x^0).\end{aligned}$$

and then substitute this in the third equation (see underlined term in the following),

$$\begin{aligned}\mathbf{x}'^{\perp} &= \mathbf{x}^{\perp} \\ \implies \mathbf{x}' - \underline{(\mathbf{x}' \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}}} &= \mathbf{x} - (\mathbf{x} \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}} \\ \implies \mathbf{x}' - \gamma (\mathbf{x} \cdot \hat{\boldsymbol{\beta}} - \beta x^0) \hat{\boldsymbol{\beta}} &= \mathbf{x} - (\mathbf{x} \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}} \\ \implies \mathbf{x}' &= \mathbf{x} + (\gamma - 1) (\mathbf{x} \cdot \hat{\boldsymbol{\beta}}) \hat{\boldsymbol{\beta}} - \gamma \boldsymbol{\beta} x^0 \\ &= \mathbf{x} + \frac{\gamma - 1}{\beta^2} (\mathbf{x} \cdot \boldsymbol{\beta}) \boldsymbol{\beta} - \gamma \boldsymbol{\beta} x^0.\end{aligned}\tag{1}$$

The first equation too can be rewritten as,

$$\begin{aligned}x'^0 &= \gamma (x^0 - \beta x^{\parallel}) \\ \implies x'^0 &= \gamma (x^0 - \mathbf{x} \cdot \boldsymbol{\beta}).\end{aligned}\tag{2}$$

B. The matrix form of (1) and (2) can be obtained from their component form. The i -th component of (1)

$$\begin{aligned}x'^i &= x^i + \frac{\gamma - 1}{\beta^2} x^j \beta^j \beta^i - \gamma \beta^i x^0 \\ &= \left(\delta_j^i + \frac{\gamma - 1}{\beta^2} \beta^i \beta^j \right) x^j - \gamma \beta^i x^0\end{aligned}$$

This gives, $\Lambda^i_0 = -\gamma \beta^i$, $\Lambda^i_j = \delta_j^i + \frac{\gamma - 1}{\beta^2} \beta^i \beta^j$. Similarly, the component form of (2) gives,

$$x'^0 = \gamma x^0 - \gamma \beta^i x^i,$$

i.e., $\Lambda^0_0 = \gamma$, and $\Lambda^0_i = -\gamma\beta^i$. Now we can write down Λ in matrix form,

$$\Lambda = \begin{pmatrix} \gamma & -\gamma\beta^1 & -\gamma\beta^2 & -\gamma\beta^3 \\ -\gamma\beta^1 & 1 + \frac{\gamma-1}{\beta^2} (\beta^1)^2 & \frac{\gamma-1}{\beta^2} \beta^1 \beta^2 & \frac{\gamma-1}{\beta^2} \beta^1 \beta^3 \\ -\gamma\beta^2 & \frac{\gamma-1}{\beta^2} \beta^2 \beta^1 & 1 + \frac{\gamma-1}{\beta^2} (\beta^2)^2 & \frac{\gamma-1}{\beta^2} \beta^2 \beta^3 \\ -\gamma\beta^3 & \frac{\gamma-1}{\beta^2} \beta^3 \beta^1 & \frac{\gamma-1}{\beta^2} \beta^3 \beta^2 & 1 + \frac{\gamma-1}{\beta^2} (\beta^3)^2 \end{pmatrix}.$$

2. From the defining equation for a Lorentz transformation,

$$\Lambda^T \eta \Lambda = \eta \tag{3}$$

show that the 00-component (top left corner entry) of this equation gives,

$$(\Lambda^0_0)^2 \geq 1,$$

i.e., either,

$$\Lambda^0_0 \geq 1,$$

or,

$$\Lambda^0_0 \leq -1.$$

Solution:

Eq. (3) in matrix form can be expressed as,

$$\Lambda^\alpha_\mu \eta_{\alpha\beta} \Lambda^\beta_\nu = \eta_{\mu\nu}.$$

The $\mu = \nu = 0$ component of this equation is,

$$\begin{aligned} \eta_{00} &= \Lambda^\alpha_0 \eta_{\alpha\beta} \Lambda^\beta_0 \\ \implies 1 &= \Lambda^0_0 \eta_{00} \Lambda^0_0 + \Lambda^i_0 \eta_{ii} \Lambda^i_0 \\ \implies 1 &= (\Lambda^0_0)^2 - \Lambda^i_0 \Lambda^i_0 \end{aligned}$$

or,

$$\begin{aligned} (\Lambda^0_0)^2 &= 1 + \Lambda^i_0 \Lambda^i_0 \\ \implies \Lambda^0_0 &= \sqrt{1 + \Lambda^i_0 \Lambda^i_0} \quad \text{or} \quad \Lambda^0_0 = -\sqrt{1 + \Lambda^i_0 \Lambda^i_0} \\ \implies \Lambda^0_0 &\geq 1 \quad \text{or} \quad \Lambda^0_0 \leq -1, \end{aligned}$$

since, $\Lambda^i_0 \Lambda^i_0 = (\Lambda^1_0)^2 + (\Lambda^2_0)^2 + (\Lambda^3_0)^2 \geq 0$ as it is a sum of squares which has to be positive semidefinite.